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The Crystal and Molecular Structure of 9-Fluorenone

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The crystal structure of 9-fluorenone, $C_{13}H_8O$, a molecule with interesting spectroscopic and photochemical properties, has been determined from counter-recorded, single-crystal X-ray diffraction data. The crystals are orthorhombic; space group *Pcab*, with $a=16\cdot068$ (4), $b=18\cdot650$ (6), $c=12\cdot550$ (4) Å and Z=16. The structure was solved by direct methods and refined by least squares. Hydrogen atoms were located and refined. The final value of the agreement index is 0.056. The molecules are planar to 0.025 Å. The C-C bonds of the benzo rings average 1.383 (1) Å, but can be divided into two sets of alternate bonds which average 1.378 (1) and 1.388 (1) Å. The angles of the benzo rings can be divided into three distinct sets which average 118.0 (2), 120.4 (2), and 121.7 (2)°, the low average being at the C(1) and C(4) positions. The C-C bonds between the benzo rings are consistent with single bonds between sp^2 carbon atoms.

Introduction

The spectroscopic properties and photochemistry of 9-fluorenone have attracted increased attention in recent years (Kuboyama, 1964; Yoshihara & Kearns, 1966; Kuroda & Kunii, 1967; Liptay, Weisenberger, Tiemann, Eberlein & Konopka, 1968; Davis, Carapellucci, Szoc & Gresser, 1969; Singer, 1969; Caldwell, 1969; Marchetti, 1971). The first X-ray investigation of 9-fluorenone was by Iball (1936), in which he established the space group and lattice constants and suggested a possible molecular packing arrangement. Griffiths & Hine (1970) have reported the crystal structure of a bromine derivative, 2-bromo-9-fluorenone. The present structure determination of 9-fluorenone was undertaken to determine more accurate bond lengths and to describe the molecular packing.

Crystal data

9-Fluorenone $C_{13}H_8O$ M.W. 180·21 Yellow, long prismatic (*c*) crystals Orthorhombic (Mo $K\alpha_1 = 0.70926$ Å) $a = 16.068 \pm 0.004, b = 18.650 \pm 0.006, c = 12.550 \pm 0.004$ Å Volume of unit cell, 3760.8 ± 2.0 Å³ Density, calculated (Z=16), 1.273 g.cm⁻³

measured (flotation), 1.24 g.cm⁻³

Linear absorption coefficient, $\mu = 0.86 \text{ cm}^{-1}$ (Mo K α) Total number of electrons per unit cell, F(000) = 1504Absent spectra: 0kl for l odd, h0l for h odd, hk0 for k odd

Space group, $Pcab (D_{2h}^{15})$

General positions: $\pm (x, y, z; \frac{1}{2} + x, \frac{1}{2} - y, z;$ $x, \frac{1}{2} + y, \frac{1}{2} - z; \frac{1}{2} + x, -y, \frac{1}{2} - z)$

Experimental

Single crystals grown from methanol were supplied by S. J. Marino of these laboratories. Space-group extinctions and preliminary lattice dimensions were obtained from precession photographs. A small crystal, $0.28 \times$ 0.35×0.47 mm, was mounted on a thin glass rod with G. E. varnish cement and covered with a glass capillary to prevent sublimation. This crystal was used for intensity data collection on an automatic Picker 4-circle goniostat. To minimize the effect of multiple reflections (Zachariasen, 1965), the *c* axis was oriented 4° from the φ axis. Twenty reflections, at moderately high Bragg angles (Mo $K\alpha_1$ radiation), were accurately centered through very narrow vertical and horizontal slits at a take-off angle of 0.5°. These observations were used as input for the computer program *PICK2* (Ibers, 1966), which refined the cell and orientation parameters by the least-squares method and generated the cards to control the automated goniostat.

Intensity data were collected with Zr-filtered Mo radiation at a 1.5° take-off angle. A scintillation detector was used, and the pulse-height analyzer was set for an approximately 90% window. All 3719 unique reflections $2\theta < 52^{\circ}$ were measured by the θ -2 θ scan technique at a 2θ scan rate of 1.0° min⁻¹. The scan range varied from 1.0° at low 2θ to 1.4° at high 2θ . Stationarycrystal stationary-counter background counts were taken for 20 sec at each end of the scan. For count rates above 13,000 cycles per sec, brass attenuators were automatically inserted in the diffracted beam. To check electronic and crystal stability during the data collection, the intensity of the 355 reflection was measured every 50th reflection. No systematic variation was observed for these standard intensities. The standard deviation of 355, calculated for 81 observations, was the same as that which one would expect on the basis of counting statistics.

The background for a reflection was approximated by a straight line between the two measured background points. The intensities were corrected for background and for Lorentz and polarization effects and were reduced to structure amplitudes, $|F_o|$. Absorption corrections were not applied, because of the low absorption coefficient and the small size and regularity of the crystal. The maximum error caused by absorption is estimated to be less than 1% of the intensity. Standard deviations $\sigma(I)$ were based on counting statistics and were corrected to $\sigma(F)$. Of the 3719 measured reflections, 2057 were less than $2\sigma(I)$ and were considered unobserved. The structure determination and refinement were based on the 1662 observed reflections.

Determination and refinement of structure

The observed structure factors were converted to normalized structure-factor magnitudes, $|E_H|$ (Hauptman & Karle, 1953), which were then phased by the multiple-solution computer program of Long (1965), modified by Tsai (1968), based on the reiterative application of Sayre's (1952) equation. An *E* map, calculated from the most consistent set of phases, yielded all 28 heavy atoms for the two molecules in the asymmetric unit. Refinement was by a block-diagonal leastsquares program (Trotter, 1965), which utilized a 3×3 block for coordinates and was diagonal in temperature factors. This program minimized $\sum w(F_o^* - F_c)^2$, where

$$F_{o}^{*} = K \left[1 + g \frac{(1 + \cos^{4} 2\theta)}{(1 + \cos^{2} 2\theta) \sin 2\theta} F_{o}^{2} \right]^{1/2} F_{o}$$

in which K is the scale factor, g is the extinction parameter (Zachariasen, 1963; Larson, 1967) and F_c is the usual calculated structure factor. Weights were defined as

$$w = \{ [qK\sigma(F_o)]^2 + (rF_o^*)^2 \}^{-1},\$$

where q and r were chosen by inspection to make the average of $\sum w(F_o^* - F_c)^2$ approximately constant for groups of increasing F_o 's (Cruickshank, 1965). The usual agreement index was defined as $R_1 = \sum |F_o^* - F_c| / \sum |F_o^*|$ and the weighted agreement index as $R_2 = [\sum w(F_o^* - F_c)^2 / \sum wF_o^*]^{1/2}$. The atomic scattering factors for carbon and oxygen were obtained from *International Tables for X-ray Crystallography* (1962) and for hydrogen from Stewart, Davidson & Simpson (1965).

Table 1. Final atomic parameters for nonhydrogen atoms

E.s.d.'s are given in parentheses. Anisotropic thermal parameters are in the form: $exp[-0.25(h^2a^{*2}B_{11}+...2klb^*c^*B_{23}]]$.

	x	у	z	B ₁₁	^B 22	B ₃₃	^B 12	^B 13	^B 23
0	0.41219(13)	0.40676(12)	0.51933(16)	7.78(12)	7.78(14)	6.00(11)	1.48(11)	-2.37(10)	-1.18(11)
c(1)	0.34430(19)	0.30441(17)	0.34444(25)	5 92 (17)	5.93(17)	6+57(17)	0.46(14)	-0.72(16)	-0.79(15)
C(2)	0.30381(19)	0.27315(17)	0.25914(28)	6.64(18)	6.17(18)	8.42(20)	-0.19(16)	-0.79(18)	-1.65(18)
C(3)	0.24890(20)	0.31184(17)	0.19882(26)	6.51(18)	7.75(18)	6.68(18)	-0.82(17)	-0.96(16)	-1.80(17)
C(4)	0.23212(18)	0.38352(16)	0.21982(24)	5.76(17)	6.58(17)	5+69(17)	-0.20(14)	-0.59(15)	-0.08(16)
C(5)	0.22267(19)	0.54729(16)	0.31575(24)	5.53(15)	5.71(15)	5.92(16)	-C.46(14)	-0.37(15)	0.98(15)
C(6)	0.23051(19)	0.61052(16)	0.37357(27)	5+93(17)	5.10(15)	8 • 58 (20)	-0.24(14)	0.33(18)	1.23(16)
C(7)	0.28439(19)	0.61530(16)	0.45877(27)	6.43(18)	4.95(15)	8.38(20)	-0.96(14)	0.94(17)	-0.97(17)
C(8)	0.33165(19)	0.55689(17)	0.49043(25)	5.89(17)	5+69(17)	6.12(16)	-0.23(14)	0.10(16)	-0+41(15)
C(9)	0.36245(18)	0.42316(16)	0.45033(22)	4.94(14)	6.73(17)	4.33(14)	-0.25(14)	-0.18(14)	-0.29(15)
C(10)	0:32847(16)	0.37558(16)	0.36566(22)	4.41(13)	5.51(15)	4.55(14)	0.15(13)	-0.11(13)	-0.24(14)
C(11)	0.27332(16)	0.41501(15)	0.30438(21)	4.57(13)	5.28'14)	4.27(14)	-0.75(13)	0.30(13)	-0.28(14)
C(12)	0.26882(17)	0.48868(15)	0.34679(22)	4.C2(13)	5.54(15)	4.37(14)	-0.46(13)	0.11(13)	0.90(13)
C(13)	0.32209(17)	0.49408(15)	0.43412(23)	4.76(14)	4.61(14)	4.70(14)	-0.17(13)	0.38(13)	-0:01(14)
0'	0.56374(13)	0.365P9(11)	0.29626(16)	7.14(11)	7.51(13)	6.20(11)	-0.76(10)	-1.78(10)	1.72(11)
C(1')	0.50260(19)	0.27187(16)	0.10678(25)	5.03(15)	5+27(17)	7.70(19)	0.01(14)	0.23(16)	0.21(15)
C(2))	0.46595(19)	0.24363(17)	0.01638(26)	5.79(18)	5.08(17)	8.84(20)	-0.57(14)	1.11(17)	-1.38(17)
C(31)	0.41633(20)	0.28671(17)	-0.04657(25)	6.75(19)	7.11(18)	6.50(18)	-C.59(16)	0.22(17)	-2.28(16)
C(41)	0.40261(18)	0.35854(17)	-0.02258(24)	5:59(17)	6.54(18)	5.27(16)	0.11(15)	-0.31(15)	-0.86(15)
C(51)	0.39719(19)	0.52049(16)	0.08107(25)	5.68(17)	5.33(15)	5.91(16)	0.53(14)	-0.16(15)	-C.26(15)
C(6')	0.40632(19)	0.58163(18)	0.14246(27)	6.91(19)	5.46(17)	7.82(20)	C•49(16)	0.64(18)	-0.04(17)
C(71)	0.45301(19)	0.58174(16)	0.23447(27)	7.22(18)	5.48(15)	7.31(18)	-1.13(16)	1.69(17)	-1.85(16)
C(81)	0.49265(19)	0.51990(15)	0.26935(24)	6.04(15)	6.06(17)	5.30(16)	-1.36(14)	0.57(16)	-0.50(15)
C(91)	0.51885(17)	0.38612(15)	0.22362(22)	4.51(13)	6.09(15)	4.75(14)	-0.98(13)	-0.08(13)	1.15(14)
C(10!)	0.48855(17)	0.34301(15)	0.13218(22)	4.25(13)	4.67(14)	4.59(14)	-0.37(13)	-0.12(14)	0.18(13)
C(11+)	0.43906(16)	0.38628(14)	0.06813(22)	4.13(12)	4.93(14)	4.27(13)	-0.22(13)	0.28(12)	-0.41(14)
C(12))	0.43667(16)	J•45884(15)	0.11481(21)	4.20(13)	4.58(14)	4.34(14)	-0.08(12)	0.31(13)	-0.09(13)
C11211	0.48301/171	0.45011(14)	0.20789/211	4.421131	4-96/141	4.36(14)	-0.86(13)	0.09(13)	-0.05(13)

All computations were performed on an IBM 360/65 computer.

Five cycles of refinement with isotropic temperature factors and two with anisotropic temperature factors reduced R_1 to 0.12. The largest structure factors were obviously affected by extinction, so the 223 largest F's were used with unit weighting to refine g. Three further cycles of anisotropic refinement were followed by the calculation of a difference Fourier, which revealed all 16 hydrogen atoms. The hydrogen atoms were assigned isotropic temperature factors, and all atomic parameters were refined for several cycles. The largest F's were again used to adjust g, after which two final cycles of atomic parameter refinement (q = 1.3, r = 0.02, g = 1.27×10^{-6} yielded the final agreement indices $R_1 =$ 0.056 and $R_2 = 0.045$ for the 1662 observed reflections. The standard deviation of an observation of unit weight was 1.13. The maximum parameter shift on the last cycle was 0.24σ . A final difference Fourier was featureless, with residual electron density between -0.21 and +0.21 e.Å⁻³, and verified the structure.

The final positional and thermal parameters, with standard deviations estimated from the least-squares process, are given in Tables 1 and 2. The average standard deviation of the positional parameters expressed in Å are 0.0029 for oxygen (range, 0.0023–0.0033), 0.0035 for carbon (range, 0.0027–0.0048), and 0.025 for hydrogen (range, 0.022–0.028). Another parameter that was refined was the scale factor, $K=0.8926 \pm 0.0011$. The final values of the observed and calculated structure factors and $\sigma(F_o)$, based on counting statistics, are given in Table 3.

	Tab	le	2.	Final	atomic	parameters	for	hydrogen	atoms
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E.s.d.'s are given in parentheses.

	x	У	z	В
н(1)	0.3859(16)	0.2773(14)	C.3897(19)	8.5(8
н(2)	0.3117(15)	0.2210(13)	0.2437(20)	8.3(8
4(3)	0.2202(15)	0.2895(14)	C.1351(19)	8+0(9
H(4)	0.1937(14)	0+4156(13)	0.1765(19)	6.7(7
H(5)	0.1827(14)	0.5448(12)	0.2510(22)	7.3(7
H(6)	0.1976(15)	0.6529(14)	0.3512(19)	7.3(7
H(7)	0.2886(14)	0.6597(12)	0.5005(19)	6.4(7
н(8)	0.3716(14)	0.5606(13)	0.5550(19)	6.8(7
H(1')	0.5376(15)	0.2426(14)	0.1536(18)	7.2(7
H(21)	0.4761(17)	0.1892(14)	-0.0024(22)	9.619
H(31)	0.3912(15)	C.2659(14)	-0.1118(18)	7.9(8
H(41)	0.3683(16)	0.3896(13)	-0.0714(21)	8.2(8
H(5!)	0.3598(16)	C.5185(13)	0.0147(21)	7.5(7
H(61)	0.3808(16)	0.6265(14)	0.1211(20)	8.0(8
H(7')	0.4559(15)	0.6243(13)	0.2762(20)	7.317
H(8!)	0.5294(17)	0.5198(14)	0.3379(21)	8.6(8

Description of structure

Fig. 1 shows the atomic labeling and the intramolecular distances and angles, uncorrected for thermal motion, for the two molecules in the asymmetric unit. The values shown in Fig. 1(c) are averages for equivalent bonds and angles assuming C_{2v} symmetry. The cell-parameter errors and the coordinate standard deviations obtained from the least-squares refinement were used to calculate estimated standard deviations for the distances and angles. These e.s.d.'s average 0.004 Å for C–O, 0.005 Å

for C–C, and 0.025 Å for C–H. The e.s.d.'s for the angles are 0.33° for C–C–O, 0.31° for C–C–C, and 1.5° for C–C–H.

The 24 C-C bond distances of the benzo rings average 1.383 Å, slightly shorter than the accepted value of 1.394 Å (Sutton, 1965). All of the bonds are within $\pm 3\sigma$ of the average, but there is an apparent systematic alternation of bond lengths around the rings. The 12 bonds of the type C(1)-C(2), C(3)-C(4), and C(10)-C(11) range from 1.382 to 1.394 Å and average 1.388 Å.



Fig. 1. Atomic labeling and bond lengths (a) and angles (b) for the two molecules in the asymmetric unit. Upper numbers are for the unprimed molecule, lower ones are for the primed molecule. Averages for equivalent bonds are given in (c).

The 12 alternating bonds of the type C(2)–C(3), C(4)– C(11) and C(1)–C(10) range from 1.368 to 1.383 Å and average 1.378 Å. The two ranges overlap only slightly and the 0.010 Å difference in the averages appears significant when compared to the calculated standard deviation of 0.001 Å for the averages. The angles of the benzo rings average the expected 120.0° but also show systematic variation. The 24 angles can be grouped into three nonoverlapping sets of eight angles each, the angles of each set being in *para* positions. The angles at C(5), C(8), and their equivalents average $118\cdot0^{\circ}$ (range, $117\cdot4-119\cdot0^{\circ}$), those at C(7) and C(12) average $120\cdot4^{\circ}$ (range $119\cdot6-121\cdot0^{\circ}$) and those at C(6) and C(13) average $121\cdot7^{\circ}$ (range, $121\cdot1-122\cdot3^{\circ}$). These differences of 2·4 and $1\cdot3^{\circ}$ appear significant when compared to calculated standard deviations of 0·2° for the averages. The variations in distances and angles are consistent with the less precise values found by Griffiths & Hine (1970) for 2-bromo-9-fluorenone, and with the SCF MO CI calculations of Kuroda & Kunii (1967) who

Table 3. Observed and calculated structure factors

Column headings are h, k, $10F_0^*$, $10|F_c|$, and $10K\sigma(F_o)$.

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predicted that the longest benzo bond would be the C(10)-C(11) bond and the shortest would be C(2)-C(3).

The relatively low values for the angles at the C(1)and C(4)-type positions are similar to those found in the photodimer of 1,4-epoxy-1,4-dihydronaphthalene (118°) (Bordner, Stanford & Dickerson, 1970), 2-acetyl-3-indazolinone (117.5°) (Smith & Barrett, 1969), 3-(pbromophenyl)phthalide (113°) (Kalyani & Vijayan, 1969), and meso-3,3'-di(p-bromophenyl)bi-3-phthalidyl (116·4°) (Kalyani, Manohar & Mani, 1967), where benzo rings are fused to five-membered rings, and in biphenylene (115.2°) (Fawcett & Trotter, 1966), 3,4:7,8-dibenzotricyclo[4·2·0·0^{2,5}]octa-3,7-diene and (115.4°) (Barnett & Davis, 1970), where benzo rings are fused to four-membered rings. For these otherwise unsubstituted benzo rings, the four- and five-membered

C(13)

-10

Q

٥

fused rings impose bond angles of approximately 90 and 108° at sp^2 carbon atoms, and cause the benzo angles at the points of fusion to become greater than 120° to relieve the strain. The effect of the enlargement of the benzo angles at the fusion points is to reduce the angles at the C(1)- and C(4)-type positions below 120° . A striking example of this effect is found in benzo-[1,2:4,5]dicyclobutene (Lawrence & MacDonald, 1969), where four-membered rings are fused to opposite sides of a benzo ring and the angle at C(1) is reduced to 108°.

The averages of four C(9)-C(13)-type bonds (1.486) Å) and of two C(11)–C(12)-type bonds (1.475 Å) are close to the 1.48 Å expected for a single bond between sp² carbon atoms (Lide, 1962; Dewar & Schmeising, 1960), and indicate little, if any, conjugation interaction

> D' 27* - 3 -2 5 - 5 - 24* -26* -9* 15* 19* 4 1

- 7*

1

Table 4. Equations of least-squares planes and distances ($Å \times 10^3$) of atoms from these planes

X, Y, and Z are expressed in Å units relative to the crystallographic axes. Planes were calculated by the method of Schomaker. Waser, Marsh & Bergman (1959) as modified by Blow (1960).

Unprimeo	d molecule	A: 0.7 B: 0.7 C: 0.7 D: 0.7	440X + 0.2 374X + 0.2 468X + 0.2 470X + 0.2	$2751 Y - 0^{\circ} 2858 Y - 0^{\circ} 2716 Y - 0^{\circ} 2685 Y - 0^{\circ} 2685 Y - 0^{\circ} 0^{\circ} $	6089Z = 3.0584 6120Z = 3.1236 6070Z = 3.0611 6082Z = 3.0317			
Primeo	i molecule	A': 0.8 B': 0.8 C': 0.8 D': 0.8	$ \begin{array}{c} 3081X + 0.2 \\ 3117X + 0.2 \\ 3082X + 0.2 \\ 3028X + 0.2 \end{array} $	2563 Y - 0.2528 Y - 0.2578 Y - 0.2578 Y - 0.2605 Y -	5304Z = 7.0983 5265Z = 7.0995 5294Z = 7.1110 5363Z = 7.0822			
	τ	Jnprime	d molecul	e		Primed	molecule	
Plane:	A	В	С	D	A'	B'	C'	
0	12	61*	10*	12*	2	-20*	-1*	
C(1)	12	68*	11*	4	-16	- 32*	- 14*	
C(2)	5	59*	6*	-4	-7	-18*	- 3*	
C(3)	2	40*	9*	1	12	10*	16*	
C(4)	-5	18*	7*	3	. 7	9*	9*	
C(5)	1	-6	22*	29*	-7	1	11*	-
C(6)	25	8	48*	60*	-9	0	-16*	-
C(7)	7	1	26*	39*	-3	-2		
C(8)	-16	-7		6*	10	3	3*	
C(9)	5	33*	-1	0*	5	- 10*	2	
C(10)	-2	39*	2	0	-5	-16*	-5	
C(11)	-13	11*	-3	-3	5	2*	5	
C(12)	-13	- 2	1	6*	0	0	-3	

6* * Not included in calculation of the plane

5

· 2



Fig.2. Stereoscopic view of the molecular packing down the b axis. The c axis is horizontal and a is vertical.

across the central ring. The C–O distances average 1.220 Å, comparable to the 1.215 Å found for aliphatic and conjugated aldehydes and ketones (Sutton, 1965). As is typical of X-ray structures, the C–H bonds average 1.01 Å, less than the true 1.08 Å separation.

Equations of least-squares planes for the molecules and for the rings of the molecules, along with deviations of the atoms from the planes, are given in Table 4. Both molecules are approximately planar with maximum deviations of 0.025 Å for the unprimed molecule (Plane A) and 0.016 Å for the primed molecule (Plane A'). The deviations from the molecular planes appear to be systematic, particularly for the unprimed molecule, but the dihedral angles between the planes of the individual rings are only about 1° and are not very significant in terms of bonding in the molecules.

The molecular packing is shown stereoscopically (Johnson, 1965) in Fig. 2. The molecules pack on end in layers perpendicular to **b**. The two molecules in the asymmetric unit are almost parallel (dihedral angle, $5 \cdot 9^{\circ}$) and average $3 \cdot 6$ Å apart. Two asymmetric units are related by a center of symmetry to form a stack of four molecules in which the two inner (primed) molecules are parallel and $3 \cdot 57$ Å apart. These stacks of four are tilted in the *ac* plane and combine with stacks related by the **c** translation to form zigzag chains. Neighboring chains are nested along **a**, but with the axis of the stacks in the opposite direction. There are no intermolecular contacts shorter than would be expected from van der Waals radii.

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